

Detection and characterization of the singularities of functions

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Francqui chair inaugural lecture

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March 28th 2024



Jean Morlet

Alex Grossmann

Ingrid Daubechies



Christine de Mol



Jean-Pierre Antoine



Christine de Mol



Jean-Pierre Antoine



Françoise Bastin



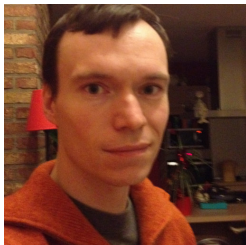
Christine de Mol



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Françoise Bastin



Samuel Nicolay



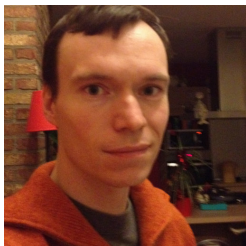
Christine de Mol



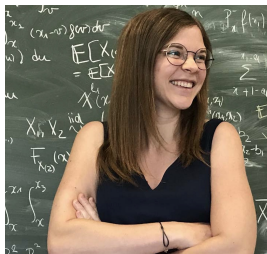
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Céline Esser



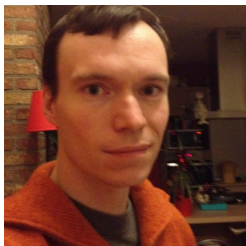
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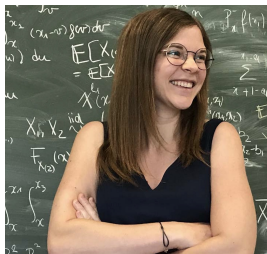
Jean-Pierre Antoine



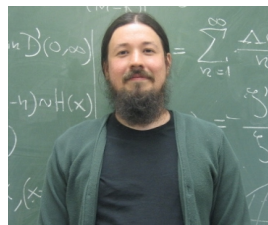
Françoise Bastin



Samuel Nicolay

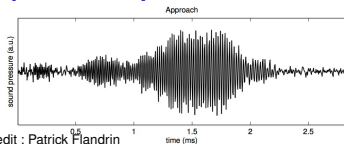


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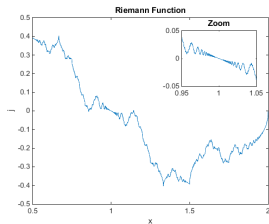
Jasson Vindas

Chirps everywhere

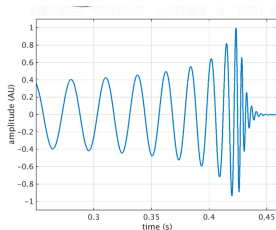


Credit : Patrick Flandrin

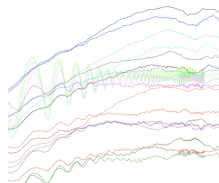
Ultrasound emitted by a bat



$$\mathcal{R}(x) = \sum_{n=1}^{\infty} \frac{\sin(\pi n^2 x)}{n^2}$$

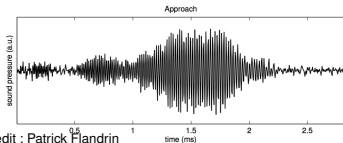


Gravitational wave



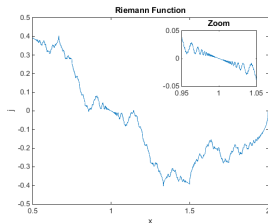
EEG

Chirps everywhere



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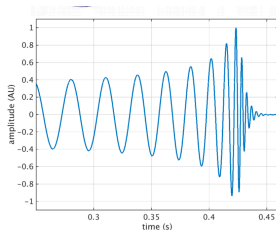


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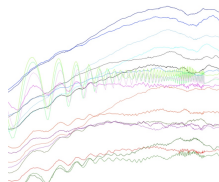
First definition : $a(t) \cos(\varphi(t))$

where $a(t)$ and $\varphi(t)$ have a “slow and smooth” evolution

AM-FM signals (*Amplitude Modulated - Frequency Modulated*)



Gravitational wave



EEG

September 14th 2015 09 :50 :45 UTC

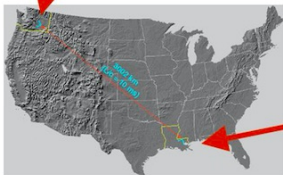
The LIGO (**L**aser **I**nterferometer **G**ravitational-Wave **O**bservatory) observatories in Hanford (state of Washington) and Livingston (Louisiana) performed the first detection of a gravitational wave

Hanford (H1=4km, H2=2km)

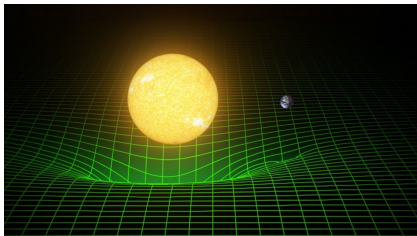
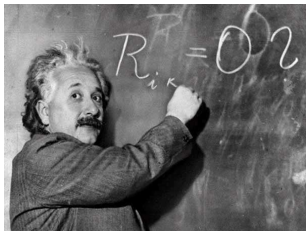


Observation of nearly simultaneous signals 3000 km apart rules out terrestrial artifacts

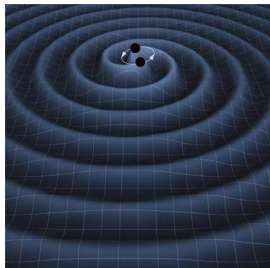
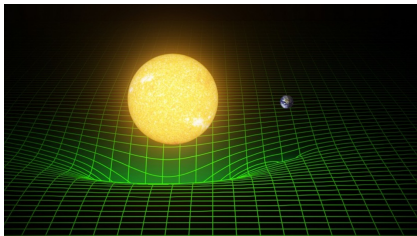
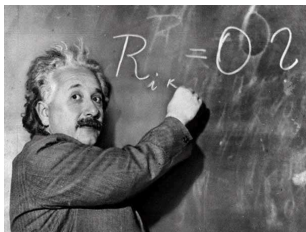
Livingston (L1=4km)



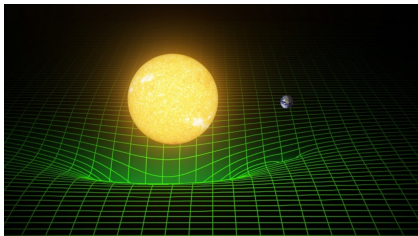
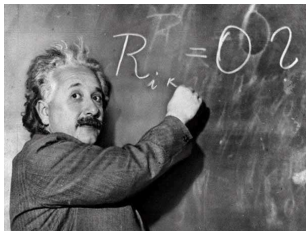
One minute of general relativity



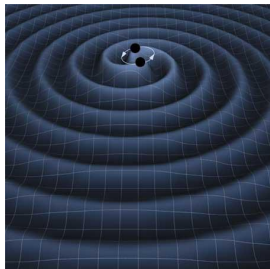
One minute of general relativity



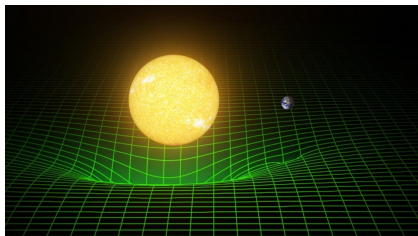
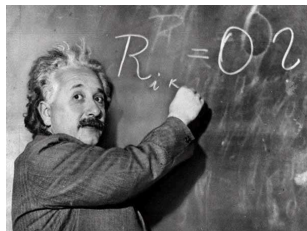
One minute of general relativity



Einstein predicted the existence of gravitational waves (1916)



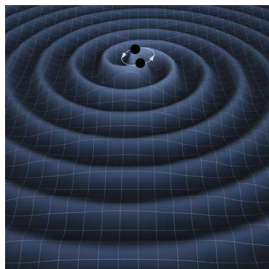
One minute of general relativity



Einstein predicted the existence of gravitational waves (1916)

Consequences of the detection :

- Confirmation of general relativity in extreme conditions of mass and energy
- A new astronomy



A few orders of magnitude

Signal emitted **1,4 billion years ago**

Coalescence of 2 black holes of **36 and 29 solar masses**

Energy dissipated in 0.2 seconds : **3 solar masses**



Credit : <http://www.gravity.phys.uwm.edu/research/highlights/index.html?artfile=160211-51.xml>

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Credit : <http://www.gravity.phys.uwm.edu/research/highlights/index.html?artfile=160211-51.xml>

Size of the recorded signal before denoising : $\sim 10^{-18}$ m

Size of the gravitational wave which crossed the earth : 10^{-21} m

Radius of the hydrogen atom : 10^{-11} m

Radius of the atomic nucleus : 10^{-15} m

Scientific challenges

Instrumental and Physics

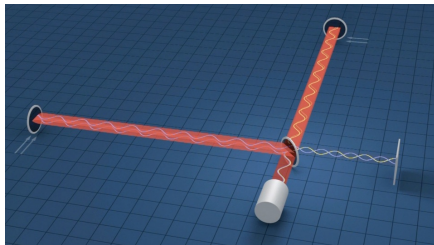
An extreme sophistication
of the experimental device

Michelson Interferometer
with 2 arms of 4 km length

The laser beam is reflected
several hundreds of times

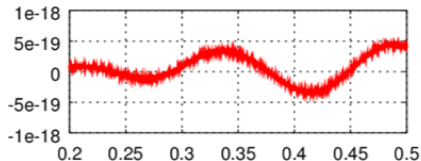


Rainer Weiss

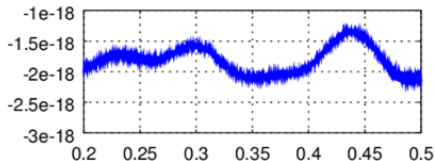


The denoising algorithm : Frequency filtering

Hanford H1: raw data



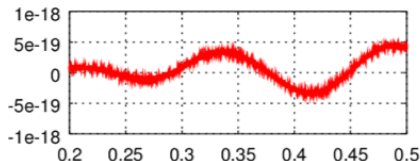
Livingston L1: raw data



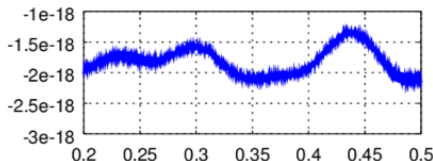
Gravitational wave GW150914 recorded by :
LIGO Hanford (H1, left) and Livingston (L1, right) detectors

The denoising algorithm : Frequency filtering

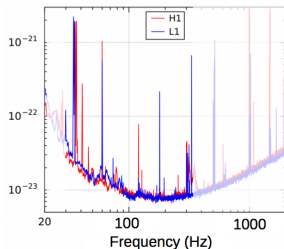
Hanford H1: raw data



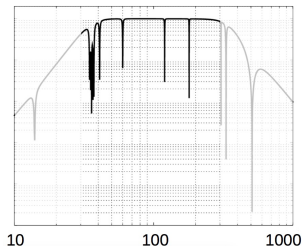
Livingston L1: raw data



Gravitational wave GW150914 recorded by :
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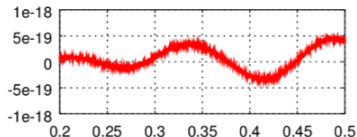
Frequencies in the data



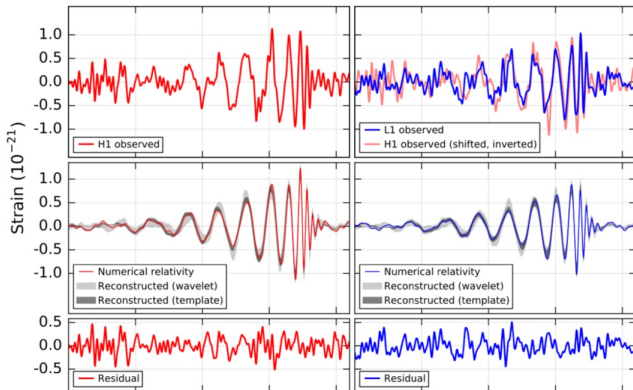
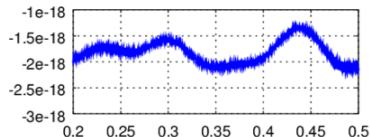
Fourier transform of the filter

The result of Fourier filtering

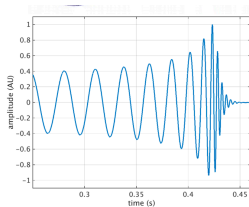
Hanford H1: raw data



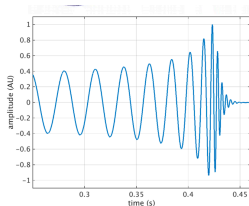
Livingston L1: raw data



The Short-Time Fourier Transform (STFT)



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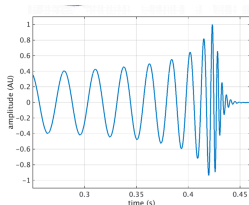


Let φ be a Gaussian function

The short-time Fourier transform of f is

$$G_f(x, \xi) = \int_{\mathbb{R}} f(t) \varphi(t - x) e^{-2i\pi t \xi} dt$$

The Short-Time Fourier Transform (STFT)

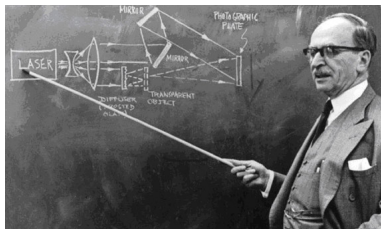


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$$f(t) = \int \int G_f(x, \xi) e^{2i\pi \xi t} \varphi(t - x) d\xi dx$$



D. Gabor

1971 Nobel prize laureate
for holography

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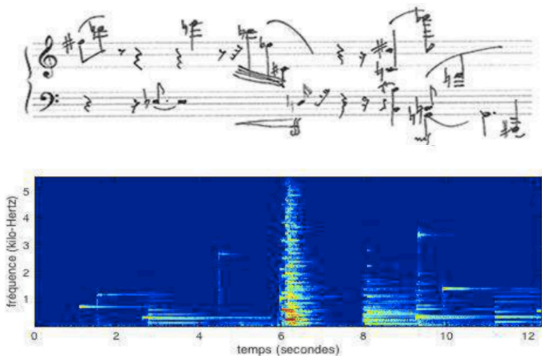
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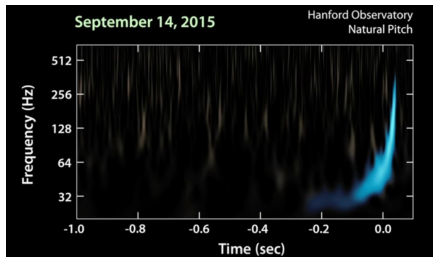
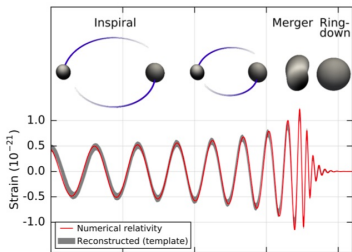
The Short-Time Fourier Transform (STFT)

$$G_f(x, \xi) = \int_{\mathbb{R}} f(x) \varphi(t - x) e^{-2i\pi t \xi} dt$$



Credit : Patrick Flandrin

STFT of gravitational waves



B. Torresani



P. Flandrin



E. Chassande-Mottin

Time-frequency orthonormal bases ?

$$(1) \quad 1_{[k, k+1[}(x) e^{2i\pi n x} \quad k, n \in \mathbb{Z}$$

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The **Balian-Low theorem** (1981) : If

$$\int (1 + t^2) |g(t)|^2 dt < \infty \quad \text{et} \quad \int (1 + \xi^2) |\hat{g}(\xi)|^2 d\xi < \infty$$

then any system of the form

$$g(x - ak) e^{ibnx} \quad k, n \in \mathbb{Z}$$

is either incomplete or over-complete, and never is a basis of $L^2(\mathbb{R})$

(1) is compatible with the Balian-Low theorem

Beyond the Balian-Low theorem

T. Steger : A Riesz basis $\varphi_n(x)$ of $L^2(\mathbb{R})$ cannot simultaneously verify
 $\exists a_n, b_n : \int (1+|t-a_n|^2)|\varphi_n(t)|^2 dt < \infty$ and $\int (1+|\xi-b_n|^2)|\widehat{\varphi_n}(\xi)|^2 d\xi < \infty$

(“strong” uncertainty principle for bases)

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How to beat Balian-Low ?

Beyond the Balian-Low theorem

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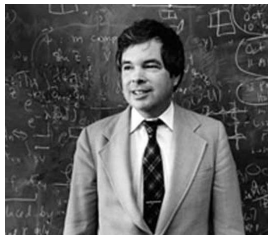
J. Bourgain constructed a basis with an optimal time-frequency localization (but far from meeting the requirements of signal processing)



How to beat Balian-Low ?

In 1987 **K. Wilson** (1982 Nobel laureate for renormalization theory) proposed a way :

Allow a double Fourier localization around two frequencies of same amplitude and opposite signs

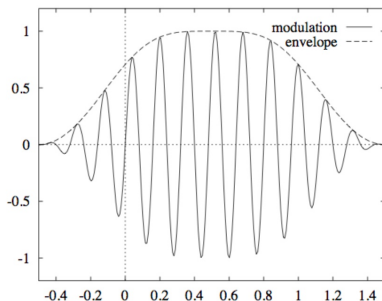


Wilson bases

Wilson bases (I. Daubechies, S. J., J.-L. Journé, 1991) are orthonormal bases of the form

$$\varphi_{0,n}(t) = \varphi(t - n) \quad n \in \mathbb{Z},$$

$$\varphi_{l,n}(t) = \begin{cases} \sqrt{2}\varphi\left(t - \frac{n}{2}\right) \cos(2\pi lt) & \text{if } l + n \in 2\mathbb{Z}, \\ \sqrt{2}\varphi\left(t - \frac{n}{2}\right) \sin(2\pi lt) & \text{if } l + n \in 2\mathbb{Z} + 1 \end{cases}$$



Wilson bases

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- ▶ φ and $\hat{\varphi}$ can both have exponential decay
- ▶ $\hat{\varphi}$ can be compactly supported



J.-L. Journé († April 2016)



I. Daubechies

Coherent Wave Burst

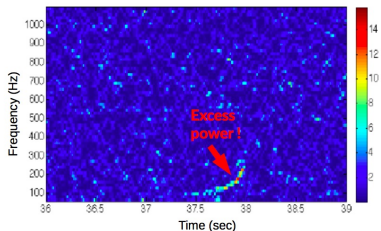
Algorithm due to S. Klimenko
and his collaborators based on a
union of dilated Wilson bases



Coherent Wave Burst

Algorithm due to S. Klimenko
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Discrete
time-frequency
analysis
of the gravitational
wave

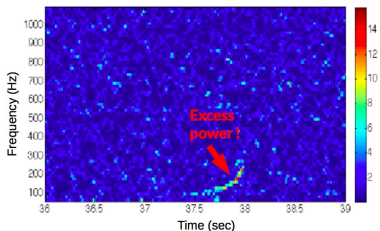


Coherent Wave Burst

Algorithm due to S. Klimenko and his collaborators based on a **union of dilated Wilson bases**



Discrete
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Advantages :

- ▶ **Sparse representation** of gravitational waves
- ▶ Fast decomposition algorithms (based on FFT)

Bottlenecks :

- ▶ Stability properties of the algorithm
- ▶ Which functions are sparse on such redundant bases ?

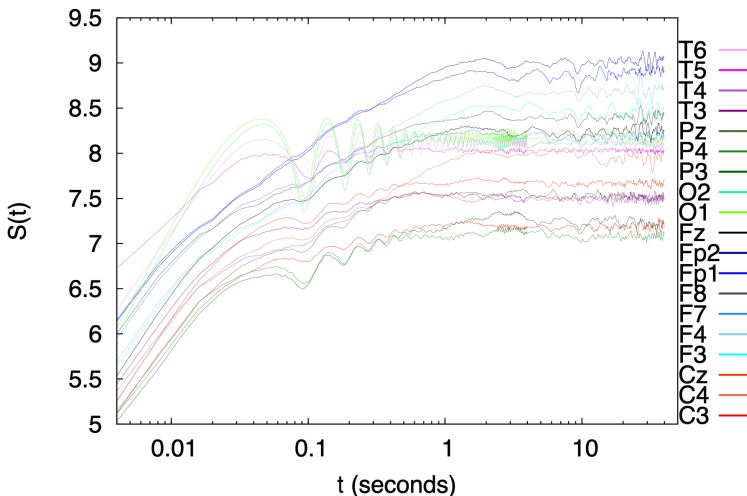
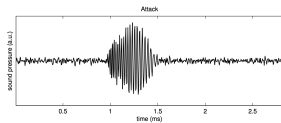
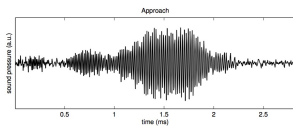


FIG. 2: The diffusion entropy $S(t)$ of the EEG increments of all the 19 channels for one of the 20 subjects considered in this study.

Credit : Dynamics of EEG Entropy : beyond signal plus noise

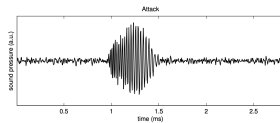
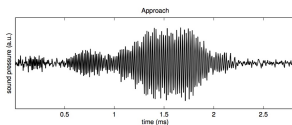
M. Ignaccolo, M. Latka, W. Jernajczyk, P. Grigolini and B.J. West
(2009)

Back to bats

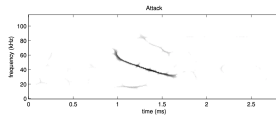
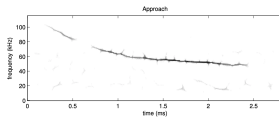


Ultrasounds emitted by a bat

Back to bats



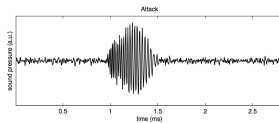
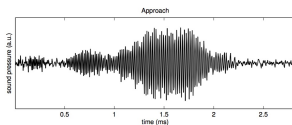
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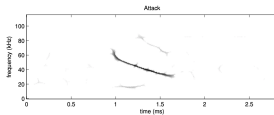
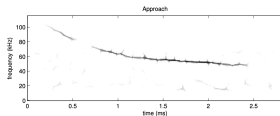
Time-Frequency analysis

Credit : *Explorations in Time-frequency Analysis* by P. Flandrin

Back to bats



Ultrasounds emitted by a bat



Time-Frequency analysis

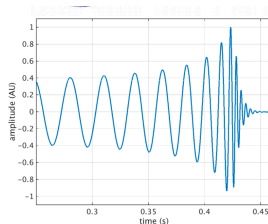
Credit : *Explorations in Time-frequency Analysis* by P. Flandrin

A more precise definition of Chirps :

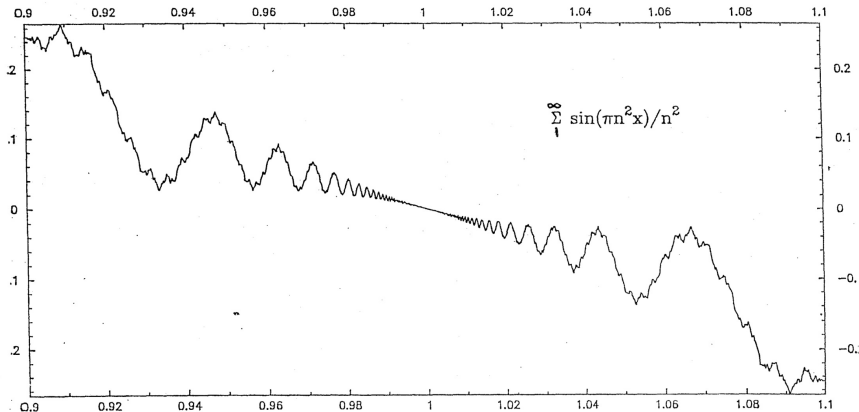
$$f(t) = a(t) \cos(\varphi(t)) \quad \text{where}$$

$$\left| \frac{a'(t)}{a(t)} \right| \ll \varphi'(t) \quad \text{and} \quad \frac{\varphi''(t)}{(\varphi'(t))^2} \ll 1$$

$$|t - t_0|^{-1/4} \cos(\omega |t - t_0|^{5/8} + \varphi)$$



Riemann's nondifferentiable function



Zoom of Riemann's function at a chirp

From G. Hardy to F. Broucke and J. Vindas (2023)

Riemann's nondifferentiable function and turbulence

A new and surprizing connexion between Riemann's function

$$\mathcal{R}(x) = \sum_{n=1}^{\infty} \frac{\sin(\pi n^2 x)}{n^2}$$

and turbulence was recently uncovered between Riemann's function and turbulence (V. Banica, A. Boritchev, D. Eceizabarrena, C. J. Garcia-Cervera, F. de la Hoz, R. L. Jerrard, D. Smets, L. Vega, and V. Vilaça Da Rocha) :

Riemann's nondifferentiable function and turbulence

A new and surprising connexion between Riemann's function

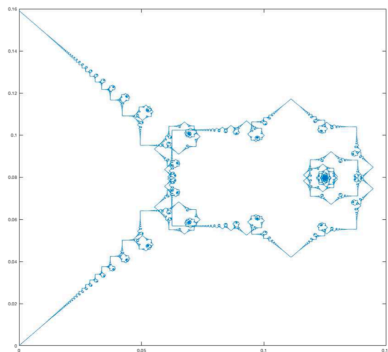
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The complex-valued version

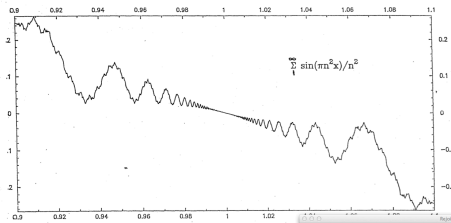
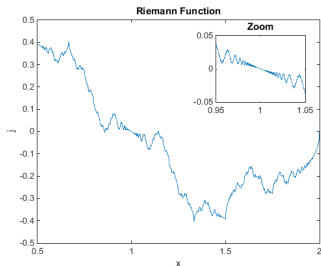
$$\phi(x) = i\pi x + \sum_{n=1}^{\infty} \frac{e^{i\pi n^2 x}}{n^2}$$

appears as the trajectory of the corners of polygonal vortex filaments that follow the binormal flow



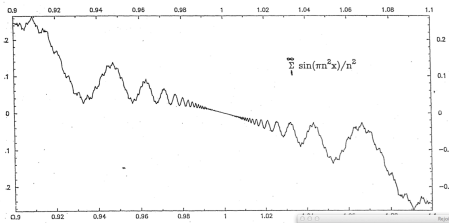
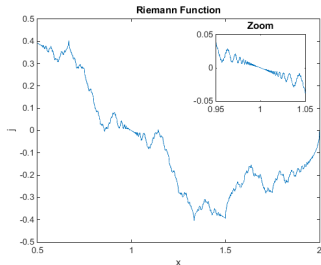
Riemann's nondifferentiable function

$$\mathcal{R}(x) = \sum_{n=1}^{\infty} \frac{\sin(\pi n^2 x)}{n^2}$$



Riemann's nondifferentiable function

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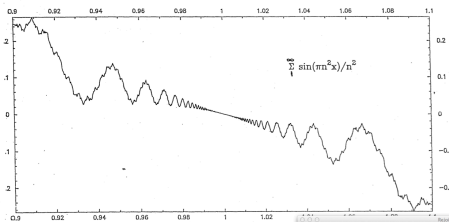
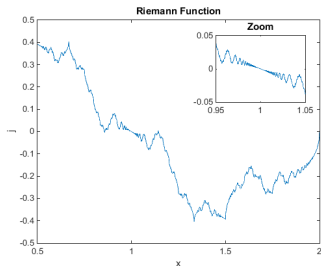
$$\mathcal{R}(\pi + x) = L(x) + x^{3/2} \left(\tilde{\mathcal{R}} \left(\frac{1}{x} \right) + x \tilde{\mathcal{R}}_1 \left(\frac{1}{x} \right) + \cdots \right),$$

where (Y. Meyer and S. J.) :

- ▶ $L(x)$ is affine,
- ▶ $\tilde{\mathcal{R}}$ is (essentially) Riemann's function,
- ▶ $\tilde{\mathcal{R}}_1$ its primitive, ...

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How can we characterize such behaviors ?

Pointwise regularity exponent

$f \in C^\alpha(x_0)$ it there exist $C > 0$ and a polynomial P of degree $< \alpha$ such that

$$|f(x) - P(x - x_0)| \leq C|x - x_0|^\alpha$$

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$$h_f(x_0) = \sup\{\alpha : f \in C^\alpha(x_0)\}$$

Pointwise regularity exponent

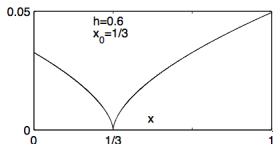
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Cusps $C_H(x) = |x - x_0|^H$



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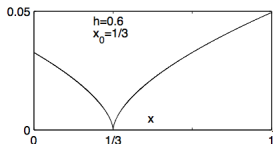
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Cusps

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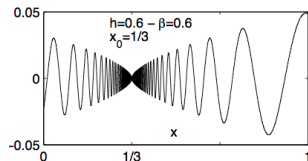
$$h_f(x) = H \quad \text{and} \quad h_{f(-)}(x) = H + 1$$



Pointwise regularity of oscillating singularities

Oscillating singularities

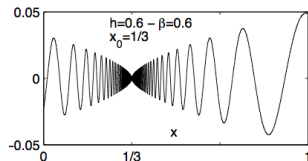
$$C_{H,\beta} = |x - x_0|^H \sin\left(\frac{1}{|x - x_0|^\beta}\right)$$



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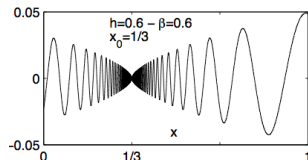
Oscillating singularities satisfy

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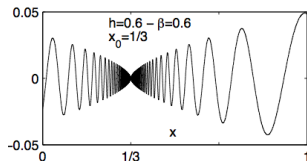
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How can one take advantage of this difference ?

Pointwise regularity of oscillating singularities

Oscillating singularities

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How can one take advantage of this difference ?

Use the wavelet characterization
of pointwise regularity

Orthonormal wavelet bases

A **wavelet basis** on \mathbb{R} is generated by one smooth well localized oscillating wavelet ψ such that the $2^{j/2}\psi(2^j x - k)$, $j, k \in \mathbb{Z}$ form an orthonormal basis of $L^2(\mathbb{R})$

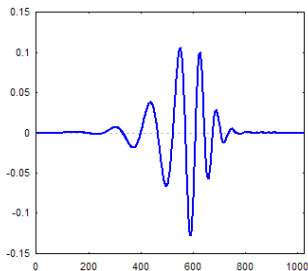
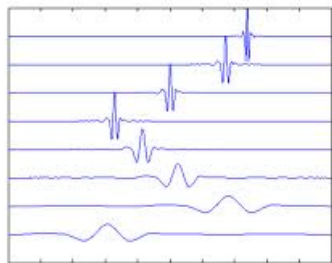
$$\forall f \in L^2(\mathbb{R}),$$

$$f(x) = \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}} c_{j,k} 2^{j/2} \psi(2^j x - k)$$

with

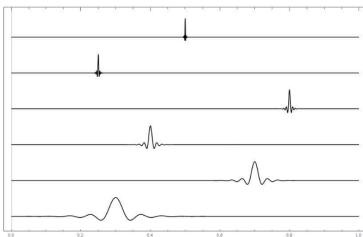
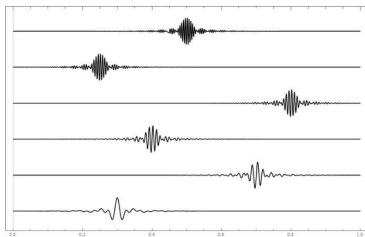
$$c_{j,k} = 2^{j/2} \int f(x) \psi(2^j x - k) dx$$

Example : $\psi = 1_{[0,1/2[} - 1_{[1/2,1[}$
(Haar wavelet)



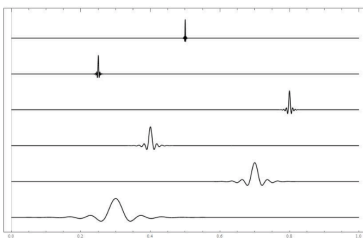
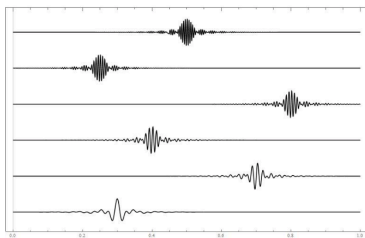
Daubechies wavelet

Orthonormal wavelet bases



Wilson vs. wavelet bases

Orthonormal wavelet bases



Wilson vs. wavelet bases

Some advantages of orthonormal wavelet bases :

- ▶ Characterization of “most” function spaces used in analysis (Hölder, Sobolev, Besov, ...)
- ▶ Characterization of pointwise regularity



Yves Meyer, Stéphane Mallat and Ingrid Daubechies

Wavelet leaders

Dyadic intervals : If $k \in \mathbb{Z}$ $\lambda = \left[\frac{k}{2^j}, \frac{k+1}{2^j} \right)$

Wavelet coefficients : $c_\lambda = 2^j \int f(x) \psi(2^j x - k) dx$

Wavelet leaders

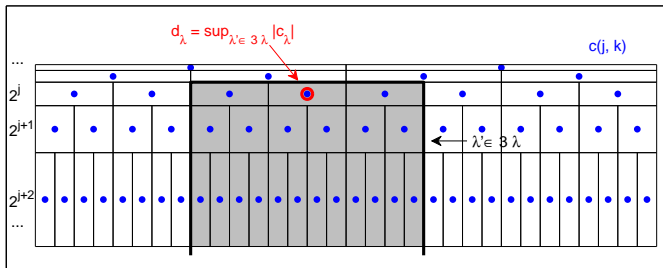
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3λ is the interval of same center as λ and three times wider

The **wavelet leaders** of a function f are the quantities

$$d_\lambda = \sup_{\lambda' \subset 3\lambda} |c_{\lambda'}|$$



Characterization of pointwise smoothness

A function f is **uniform Hölder** if $\exists \varepsilon > 0 : f \in C^\varepsilon(\mathbb{R})$

or, equivalently, if $\exists \varepsilon > 0 : \sup_{\lambda \in \Lambda_j} \leq C 2^{-\varepsilon j}$

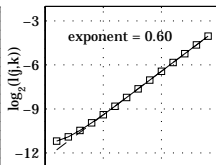
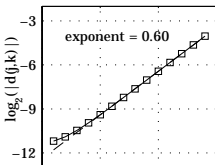
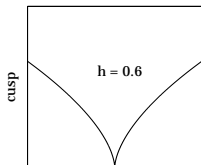
Let $\lambda_j(x_0)$ denote the dyadic cube of width 2^{-j} that contains x_0

Theorem : (S. J.) If $f \in C^\varepsilon(\mathbb{R})$ for an $\varepsilon > 0$, then

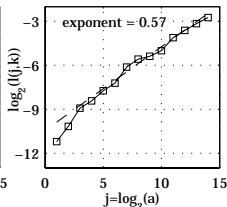
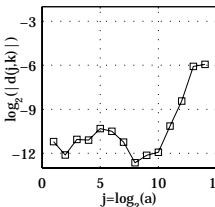
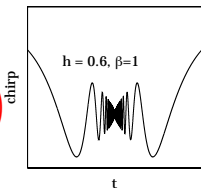
$$\forall x_0 \in \mathbb{R}^d \quad h_f(x_0) = \liminf_{j \rightarrow +\infty} \frac{\log(d_{\lambda_j(x_0)})}{\log(2^{-j})}.$$

Wavelet characterization of pointwise smoothness

$$|x - x_0|^h$$



$$|x - x_0|^h \sin\left(\frac{1}{|x - x_0|^\beta}\right)$$



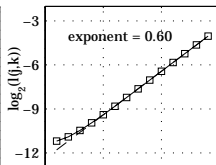
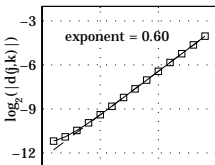
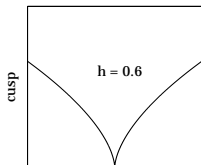
Function

Regression
on wavelet
coefficients

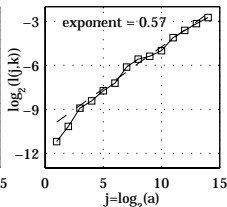
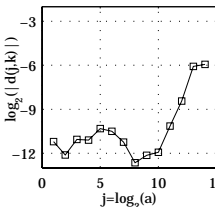
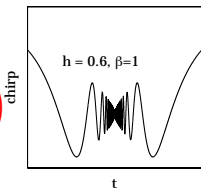
Regression
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leaders

Wavelet characterization of pointwise smoothness

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Function

Regression
on wavelet
coefficients

Regression
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leaders

Log-log regressions of wavelet leaders yield a correct estimations of pointwise Hölder exponents for all kinds of singularities

Oscillating singularities

Mathematical examples

- ▶ Riemann series : Y. Meyer and S. J.
- ▶ Random wavelet series : P. Abry, R. Leonarduzzi, C. Melot, H. Wendt, S.J. and C. Esser, B. Vedel, S.J.
- ▶ Sample paths of some specific Lévy processes : P. Balança

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Multifractal analysis allows to derive information on the size of the sets of points x where

$$h_f(x) = H \quad \text{and} \quad h_{f(-1)}(x) = H + \beta + 1$$

Next challenges in applications

- ▶ Turbulence :
with Patrice Abry (ENS Lyon)
and Françoise Argoul
(Bordeaux University)



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- ▶ Build a common framework for the two approaches

If you want to know more on these topics
you are welcome to the next lectures
and the workshop !
(April 16-19)

